

# Rupture directivity and source-process time of the September 20, 1999 Chi-Chi, Taiwan, earthquake estimated from Rayleigh-wave phase velocity

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Rayleigh-wave phase-delay differences between the main shock and its aftershocks, regarded as source duration varying with azimuth, at period range of 60 ~ 120 seconds, were used to rapidly determine the rupture directivity and the source-process time of the 1999 Chi-Chi, Taiwan, earthquake. Results show that the earthquake faulting exhibits an obvious directivity with an optimal rupture azimuth of about 42° in the northeast direction. The earthquake has an average source-process time at  $40.6 \pm 0.5$  sec and a rupture length of  $77.0 \pm 3.6$  km during rupture. In addition, periods at nodes of amplitude spectra of the main shock were adopted to estimate the rupture time, to be about  $34.0 \pm 0.7$  sec; the rise time of  $6.7 \pm 0.6$  sec, and a slightly slow rupture velocity,  $2.27 \pm 0.15$  km/sec, were also obtained. According to Savage's suggestion, the rupture width of the fault was estimated to be  $30.5 \pm 4.5$  km. Relying on these estimated faults parameters, we calculate the average static stress drop of about 56 bars, lying in between the stress drop of interplate earthquakes and intraplate earthquakes. We also infer the average dynamic stress drop of 52 bars based on Brune's theory. By the comparison of the two stress-drops, an Orowan's stress model or a stress model of frictional overshoot is probably appropriate to describe the rupture behavior of the earthquake. The radiated seismic energy, then, was estimated at about  $1.9 \times 10^{23}$  ergs, about 16 times larger than the value reported by the USGS.

## 1. Introduction

The Chi-Chi, Taiwan, earthquake, a shallow and low-angle thrust-faulting earthquake with  $M_s = 7.6$  (USGS), occurred in central Taiwan on September 20, 1999. This was the most seriously disastrous earthquake for Taiwan within fifty years. The obvious surface rupture (~80 km) and large crust deformation (vertical displacement of 1 ~ 8 m) were generated during earthquake rupture (Ma *et al.*, 1999; Wang *et al.*, 2000). In accordance with source rupture inversion (Lee and Ma, 2000; Yagi and Kikuchi, 2001), the Chi-Chi earthquake ruptures from south to north and consists of several subevents with the largest one in the northern portion of the fault. Such an earthquake would cause asymmetrically azimuthal distribution of waveforms and have the observed seismic-wave travel-time delayed because of source finiteness.

Ben-Menahem (1961) first proposed a theory of finite moving source to account for the effect of rupture propagation on far-field seismograms. The effects on seismic waves will be observed in two ways: one is to produce the time delay in a certain direction of wave propagation, and the other is that amplitude spectra are weakened and then many nodes are generated (e.g. Ben-Menahem, 1961; Kanamori and Anderson, 1975). This theory has been widely applied to determine experimentally the fault length and the rupture velocity for large and great earthquakes (e.g. Ben-Menahem,

1961; Press *et al.*, 1961; Filson and McEvelly, 1967). In recent years, the empirical Green's function analysis has been employed to investigate the source-process time and the rupture azimuth by means of surface wave data (e.g. Velasco *et al.*, 1994). Furthermore, the experimental determination of the fault parameters is likely to cause an ambiguous estimation in the rupture length or the rupture velocity on account of the trade-off relation between these fault parameters. On the other hand, the empirical Green's function analysis has strict constraints including that focal mechanism and source depth of the main shock must be analogous to that of the aftershock, treated as an empirical Green's function; in addition, the locations of the two earthquakes must be quite close each other.

In this study, we calculate the phase-delay differences of fundamental-mode Rayleigh-wave, taken as the source duration varying with azimuth, between the main shock and its aftershocks so that the inconsistent locations and focal mechanisms between the main shock and its aftershocks are capable of being removed. Accordingly, following Ben-Menahem's theory, we can estimate rapidly and objectively the rupture directivity and fault parameters of the 1999 Chi-Chi, Taiwan, earthquake by using a least-squares technique.

## 2. Data

Long-period seismograms recorded at GSN (Global Seismic Network) stations with epicentral distances between 30° ~ 90° had been used in this study. Rayleigh wave trains were drawn out from the vertical-component seismograms

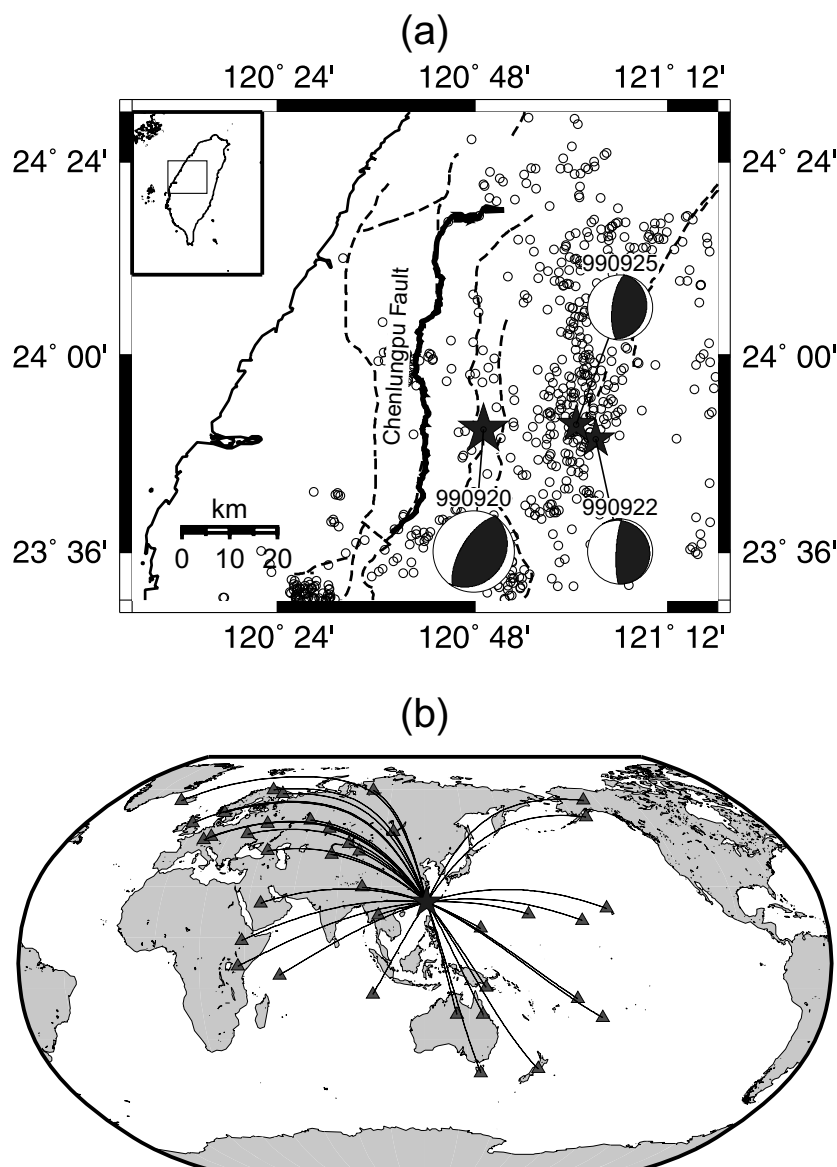


Fig. 1. (a) The large star represents the epicenter of the 1999 Chi-Chi, Taiwan, earthquake. Two small stars display the aftershocks, occurring in 09/22/1999 and 09/25/1999, respectively. Thick line denotes the surface rupture of the Chenlungpu fault, on which the main shock occurred. Dotted lines stand for the active faults in this region. Focal mechanisms, reported by the Harvard CMT group, are also shown. Aftershocks with  $M_L \geq 3.0$  also plotted in the figure by open circles. (b) Map of the GSN stations (solid triangles) used in this study.

Table 1. Source parameters used in this study.

No.	Date	Origin time (UT)			Location		Depth (km)	Ms	Fault plane solution		
		hr.	min.	sec					Strike	dip	slip
1	1999 Sept. 20	17	47	15.9	120.82°E	23.85°N	8	7.7	37°	25°	96°
2	1999 Sept. 22	00	14	40.8	121.05°E	23.83°N	26	6.4	327°	12°	55°
3	1999 Sept. 25	23	52	49.5	121.01°E	23.86°N	17	6.4	12°	20°	95°

\*Both origin time and location are provided by the Central Weather Bureau (CWB), Taiwan. Fault plane solutions are published by the Harvard CMT group. The depth of the main shock is from the final report of the CWB. The depths of the two aftershocks are from the report of the Harvard CMT group.

with a group velocity range of 2.3 ~ 4.8 km/sec. Each seismogram was corrected for the instrumental response, filtered between 10 ~ 300 seconds and tapered by a cosine function. For periods less than 60 sec, the effect of source

finiteness from the main shock is likely to distort seismograms, and result in incorrect phase-velocity measurements. At periods greater than 120 sec, Rayleigh wave trains are interfered by noise with ease. For this reason, taking the

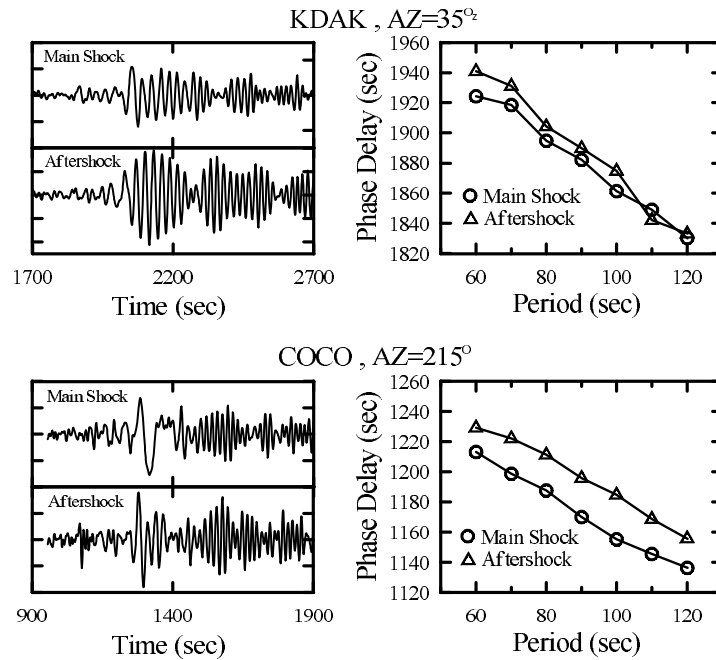


Fig. 2. Diagrams showing the vertical-component seismograms and the Rayleigh-wave phase delays between the main shock and 09/22/1999 aftershock, observed at the stations KDAK and COCO for periods of 60 ~ 120 seconds, respectively.

influence of the finite source and noise on seismic waves into consideration, we only adopted the observed data with period range of 60 ~ 120 seconds to complete the phase velocity calculations. Figure 1(a) shows the locations and focal mechanisms of the main shock and its two large aftershocks. The source parameters of these earthquakes are listed in Table 1. As shown in Fig. 1(b), the stations used in this study reveal uniform azimuthal coverage around the epicenter. In Fig. 2, the seismograms of the main shock present the larger wave duration as compared with those of 09/22/1999 aftershock, especially for records at the station COCO, which departs from the rupture direction near to 180°. The phase-delay differences of Rayleigh waves observed at the station KDAK, which approaches the rupture direction, have the small deviations, whereas in opposite direction the phase-delay differences observed at the station COCO show larger values (Fig. 2). Such an obvious directivity provides us a good opportunity to study the rupture features for this earthquake.

### 3. Method

According to the theory of Ben-Menahem (1961), for an earthquake with a unilateral faulting, the phase delay due to the rupture propagation of source is defined as  $\frac{1}{2}(\frac{L}{V_r} - \frac{L}{C} \cos \Theta)$ , where  $L$  is the fault length,  $V_r$  the rupture velocity and  $C$  the average phase velocity in the source region. The azimuth,  $\Theta$ , is measured clockwise from the rupture direction of the fault to the station. Moreover, the dislocation source time function, often considered to be a ramp function with rise time,  $\tau$ , for far-field observations, will also make the observed phase velocity slow down. Thus, the source-process time,  $T_{\text{SPT}}$ , observed at a certain azimuth (corresponding to a given station), can be written in the following form.

$$\begin{aligned} T_{\text{SPT}} &= \left( \frac{L}{V_r} - \frac{L}{C} \cos \Theta \right) + \tau \\ &= \left( \frac{L}{V_r} + \tau \right) - \frac{L}{C} \cos \Theta. \end{aligned} \quad (1)$$

In Eq. (1), there is a linear relationship between  $T_{\text{SPT}}$  and  $\cos \Theta$ . On azimuthal average, the whole source-process time is regarded as  $(L/V_r + \tau)$ , that is, the summation of the rupture time,  $L/V_r$ , and the rise time of source.

In fact, the source-process time for a given azimuth (or station) can be easily determined from the phase-delay differences of the Rayleigh-wave between the main shock and its aftershocks after some corrections have been made. In the following analysis, the phase velocities with periods of 60 ~ 120 sec are first measured for the main shock and its two aftershocks by the single-station method (cf. Yu and Mitchell, 1979) after taking away initial phase of source (cf. Wang, 1981). In addition, the deviations in estimating the source-process time should be corrected owing to the source duration of the aftershocks and the inconsistent locations between the main shock and its aftershocks. Subsequently, the Rayleigh-wave phase-delay differences,  $T_{\text{PD}}$ , corresponding to the source-process time, at a given station can be expressed as

$$T_{\text{PD}} = 2(t_m - t_a) + S_a + 2(\Delta d/C) \quad (2)$$

where  $t_m$  and  $t_a$  are the average phase delays of Rayleigh-wave of the main shock and its aftershock, respectively.  $S_a$ , the source duration of aftershock, is taken as 8.0 sec for earthquakes with  $M_s = 6.4$  when the effects of the finite source on aftershocks are neglected, and  $\Delta d$  is the difference of epicentral distance in km between the two earthquakes.  $C$  is the average phase velocity in the source region, and plays a key role in estimating the rupture length

as indicated in  $L/C$  of Eq. (1). Because the dispersion in phase velocity is likely to obscure the estimation of the rupture length, we only adopted the phase velocities at period range of 60 ~ 120 sec, where the values are almost a constant (about 4.05 km/sec) for the continental structure. From Eqs. (1) and (2),  $L/C$  and  $(L/V_r + \tau)$  can be easily estimated by a least-squares fitting.

#### 4. Rupture Directivity, Source-Process Time and Rupture Length

In Fig. 3, open circles and open triangles represent the average Rayleigh-wave phase-delay differences of the main shock to the 09/22/1999 and 09/25/1999 aftershocks, respectively. The rupture directivity is clearly exhibited with azimuth from Fig. 3, that is, there is the minimum time delay when  $\Theta = 0^\circ$ ; however, the maximum time delay appears at  $\Theta = 180^\circ$ . The inset of Fig. 3 displays the optimal rupture azimuth of about  $42^\circ$  in the northeast direction. This corresponds with the studies of the source rupture inversion (e.g. Yagi and Kikuchi, 2001) and the multiple events analysis (Kao and Chen, 2000). The best linear relation between  $T_{\text{SPT}}$  and  $\cos \Theta$  is  $T_{\text{SPT}} = (40.6 \pm 0.5) - (19.0 \pm 0.9) \cos \Theta$  by using a least-squares method. From the slope of the regression line, the fault length is estimated to be about  $77.0 \pm 3.6$  km when the average Rayleigh-wave phase velocity of 4.05 km/sec in the source region is employed. This result agrees well with field observations (Ma *et al.*, 1999; Wang *et al.*, 2000). The estimated source-process time,  $40.6 \pm 0.5$  sec, for the Chi-Chi, Taiwan earthquake is in agreement with observations of the surface ground motions (Huang, 2000), and larger than that for similar-sized earthquakes, which have been routinely reported in several earthquake catalogs. This seems to suggest that the Chi-Chi, Taiwan, earthquake is likely to have a slow rupture behavior, on average, relative to similar-sized earthquakes. In general, rupture velocity is estimated from the whole source-process time and the rup-

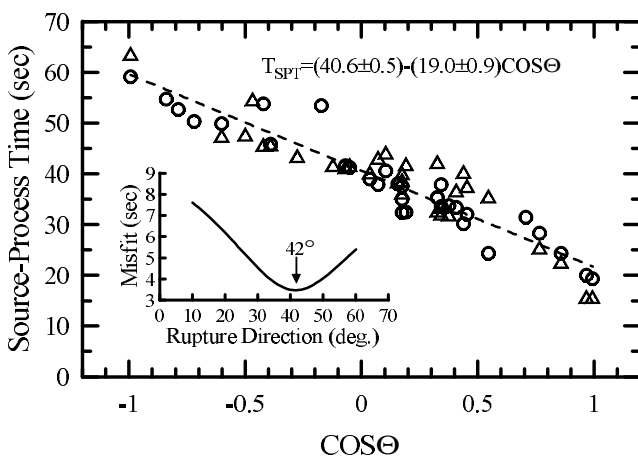


Fig. 3. The inset denotes the optimal rupture azimuth of  $42^\circ$  determined from a minimum misfit. The plot of source-process time vs.  $\cos \Theta$  presents a best linear relation based on the optimal rupture azimuth. Open circles show the phase-delay differences between the main shock and the 09/22/1999 aftershock; however, open triangles show the phase-delay differences between the main shock and the 09/25/1999 aftershock. The fitted equation is shown within the figure.

ture length when the rise time of the earthquake is ignored. Hence, the rupture velocity is about 1.9 km/sec for the Chi-Chi, Taiwan, earthquake; this value is a slightly lower than that for shallow earthquakes occurring at the other places. However, ignoring the effect of the rise time is likely to underestimate the rupture velocity.

#### 5. Rupture Velocity, Rise Time and Rupture Width

In order to obtain more reasonable rupture velocity, we must divide the whole source process-time into the rupture time and the rise time (see Eq. (1)). It is difficult, however, to determine individually the rise time from Eq. (1). Hence, an additional viewpoint is necessary in this study for objectively deducing the rupture time. In frequency-domain, the source finiteness and the rise time can be expressed in terms of a sinc function, which make many nodes in amplitude spectra. Thus, the period of the first node generated by the source finiteness corresponds to  $(\frac{L}{V_r} - \frac{L}{C} \cos \Theta)$ , where notations are given in the previous section. From global observations, the rise time is generally less than 10 sec for earthquakes of such magnitude-sized Chi-Chi earthquake (Kanamori and Anderson, 1975; Geller, 1976). In other words, at periods lower than 10 sec, the period of the first node from the source finiteness is usually contaminated by the rise time. Accordingly, in this study, we estimated the rupture time by using the first node of the amplitude spectra at periods larger than 10 seconds. Two examples of

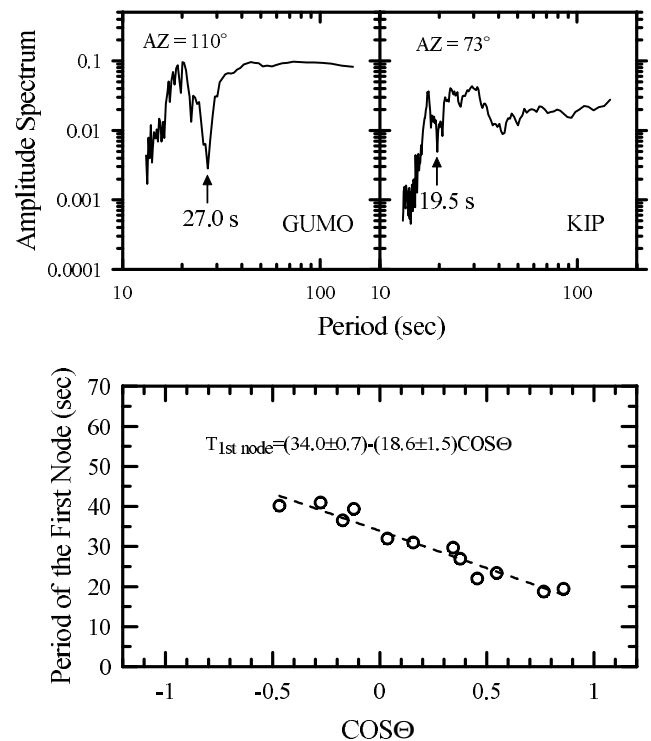


Fig. 4. Upper diagram shows two examples of amplitude spectra for the stations GUMO and KIP. Arrows designate the position of first node, and the corresponding periods are shown below. Lower diagram indicates the plot of period of the first node vs.  $\cos \Theta$ . A best linear relation from 12 data points is shown by a dotted line according to the optimal rupture azimuth (see Fig. 3). The fitted equation is also shown.

the amplitude spectra with clear nodes are shown in Fig. 4. Although only 12 data points were used, the result shows a linear relation between the period of the first node and  $\cos \Theta$ . The corresponding directivity is also shown in lower diagram of Fig. 4. Because the slope of regression line for the lower diagram of Fig. 4 is in good agreement with that for Fig. 3, the consistency between the two figures leads to reliable determinations of the rupture time and the rupture velocity. From Fig. 4, the rupture time is estimated to be  $34.0 \pm 0.7$  sec. For this reason, the rupture velocity and the rise time can be determined simultaneously as  $2.27 \pm 0.15$  km/sec and  $6.7 \pm 0.6$  sec for the earthquake. The rupture width is also estimated to be  $30.5 \pm 4.5$  km according as Savage (1972) suggested that  $\tau_r = W/2V_r$ , where  $\tau_r$  is the rise time,  $W$  the fault width and  $V_r$  the rupture velocity.

## 6. Discussion and Conclusion

The estimated fault parameters listed in Table 2 show that the Chi-Chi, Taiwan, earthquake manifests a longer source-process time as compared with those similar-sized earthquakes. The average rupture velocity is estimated at about 0.76 times the shear wave velocity (about 3.0 km/sec in the upper crust), and is slightly lower, on average, than the common value for shallow earthquakes. Hence, the rupture length and the rupture velocity are the two main factors that control the source-process time during faulting and produce obvious rupture directivity. The estimated source-process time of about 41 sec in this study agrees with Huang's observations (Huang, 2000). He used strong-motion data observed in Taiwan to reconstruct time-dependent spatial ground motions of the 1999 Chi-Chi, Taiwan, earthquake, and indicated that seismic energy radiation from the source rupture was entirely released within 42 sec after the start of the rupture.

From Fig. 3, the estimated rupture direction was  $42^\circ$  toward the northeast. In this study, we also determined the rupture direction only using the main shock but not aftershocks corrections. In Fig. 4, the nodes are purely picked from the spectra of the main shock, and then used to estimate the rupture direction as in Fig. 3. The rupture direction estimated from Fig. 4 is  $38^\circ$  toward the northeast, which is consistent with the estimate from Fig. 3. This result indicates that the fault ruptures along the fault plane from a shallower depth at the southern segment of the fault to a deeper depth at the northern portion of the fault. The rupture direction on the fault plane is finally projected onto the earth surface and result in a value of  $42^\circ$  measured clockwise from the north.

Table 2. Fault parameters estimated in this study.

Rupture length of earthquake fault	$77.0 \pm 3.6$ km
Rupture width of earthquake fault	$30.5 \pm 4.5$ km
Average source-process time	$40.6 \pm 0.5$ sec
Rupture azimuth	$42^\circ$
Rupture time	$34.0 \pm 0.7$ sec
Rise time	$6.7 \pm 0.6$ sec
Rupture velocity	$2.27 \pm 0.15$ km/sec

Although this result is different from the other source models showing a north-south direction rupture (Lee and Ma, 2000; Ma *et al.*, 2000), our result seems to be in agreement with the recent study of Yagi and Kikuchi (2001). Yagi and Kikuchi (2001) adopted the source duration of about 40 sec to invert the slip distribution of the earthquake by joint inversion of strong-motion data and teleseismic data and shown an extensive and a deeper rupture along the fault plane, especially on the northern part of the fault. The rupture depths vary from about 20 km at the southern portion of the fault to about 50 km at the northern one along the fault plane. Discrepancies between these source models are mainly from the usage of different source duration and inversion methods. A shorter source-duration used in the study of source rupture process might cause the slip distribution concentrated on the shallow depth. Huang (2000) reported that the whole energy release of the earthquake is about 42 sec, consistent with our study. Hence, we consider that the longer source duration used to invert the source rupture is reasonable. In a recent study, Huang and Wang (2001) stated that the scaling of power spectra of near-field seismograms of the earthquake increasingly vary from south to north. This seems to imply that a deeper rupture would take place at the northern part of the fault than at the southern one, and then the whole rupture with the northeast direction during faulting would be probably generated. The surface breaks of the Chelungpu fault, on which the Chi-Chi, Taiwan, earthquake occurred, almost exhibit north-south orientation, but the geometry of the fault under the ground is not known. Recently, Ji *et al.* (2001) used a model with multi-fault segments instead of a single fault model to invert the slip distribution of the Chi-Chi earthquake, and concluded that all contributions from the multi-fault portions would result in a northeast rupture direction, incompatible with the orientation of the Chelungpu fault. Consequently, our results apparently represent the whole feature of the earthquake rupture from surface-wave observations.

A rise time of  $6.7 \pm 0.6$  sec for the Chi-Chi earthquake is larger than that for earthquakes with the same magnitude (Kanamori and Anderson, 1975; Geller, 1976). This result is comparable to a recent study of Huang *et al.* (2000), in which the rise time of the northern rupture of the fault was estimated to be 5 sec. The rupture length is about 2.5 times the rupture width, which slightly departs from the geometrical similarity of the fault plane. The rupture length and rupture width of a large earthquake are often determined by the distributions of aftershocks, but for the Chi-Chi earthquake, the aftershocks are almost located on the east of the fault so that it is difficult to estimate the rupture length and rupture width (Fig. 1(a)). Therefore, the main rupture area estimated by aftershocks would probably result in overestimation and then presumably give a lower stress drop (e.g. Gibowicz, 1986).

An average dislocation of  $3.5 \pm 0.7$  meters is estimated from the seismic moment  $M_0 = 2.4 \times 10^{27}$  dyne-cm (USGS), the rigidity,  $\mu = 3.0 \times 10^{11}$  dyne-cm<sup>-2</sup> and the area of the fault,  $A$ , in light of  $M_0 = \mu AD$  (Aki, 1966). The area of the fault can be determined by the estimated rupture length and rupture width listed in Table 2. Thus the average static stress drop,  $\Delta\sigma$ , is inferred as 56 bars,

which is different from the report of EIC Seismological Note in Japan ( $\sim 33$  bars). This value seems to lie in between values for interplate earthquakes ( $\sim 30$  bars) and intraplate earthquakes ( $\sim 100$  bars) (Kanamori and Anderson, 1975). This might be possibly related to the particular tectonic feature in Taiwan. The average particle velocity of the earthquake is about 52 cm/sec, which is near common observations for large earthquakes (Kanamori, 1994). Following Brune's theory (Brune, 1970), we determine the average dynamic stress drop,  $\Delta\sigma_d$ , to be about 52 bars, which is almost the same as the average static stress drop. As a result, the Orowan's model (Orowan, 1960), in which the dynamic stress is equal to the final stress, is likely to describe the rupture behavior of the earthquake. However, since the average dynamic stress drop is somewhat lower than the average stress drop, the stress model of frictional overshoot, in which the dynamic stress is smaller than the final stress, provides an alternative interpretation to the rupture of the Chelungpu fault, on which the Chi-Chi, Taiwan, earthquake occurred. By systematically calculating the source parameters of the earthquake, an estimate of the radiated seismic energy,  $E_S$ , according to the formula:  $E_S = M_O(2\Delta\sigma_d - \Delta\sigma)/2\mu$  (Kanamori and Heaton, 2000), can be calculated. Then, the radiated seismic energy is calculated to be  $1.9 \times 10^{23}$  ergs, compared with  $1.0 \times 10^{24}$  ergs estimated by Mori *et al.* (2000) using near-field seismograms and about 16 times the value given by the USGS ( $\sim 1.2 \times 10^{22}$  ergs). Moreover, the  $E_S/M_O$  ratio (about  $8.0 \times 10^{-5}$ ) is also quite close to the global average observations ( $\sim 5.0 \times 10^{-5}$ ) (Vassiliou and Kanamori, 1982). Through such this work, we can rapidly estimate the fault parameters for a unilateral faulting earthquake from phase-delay differences of Rayleigh waves.

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